

25 **1. Introduction**

26 In recent years, additive manufacturing technologies are changing production processes in many
27 sectors such as: biomedical [1–3], metallurgical [4–6], aerospace etc. Some of the most promising
28 processes are Selective Laser Sintering (SLS) and Selective Laser Melting (SLM), in which a
29 controlled laser beam is used for sintering/melting of powders. In particular, three dimensional objects
30 can be obtained by means of a layer by layer additive technique. Consequently, significant advantages
31 of these additive manufacturing technologies with respect to traditional sintering processes [7–9] are
32 given by the possibility of advanced product customization and of complex internal morphology of
33 the artefacts [10]. The discrimination between SLS and SLM is labile and it is mostly related to the
34 amount of energy released to the material by the laser beam during the sintering/melting process [11].
35 An advantage of SLS versus SLM lies in smaller contraction of the processed material due to the
36 limited change in the particulate solids. Therefore, SLS is a more promising technique for the
37 production of high accuracy sized objects. Furthermore, SLS allows the production of artefacts with
38 internal porosity, a valuable property that is required in some applications, such as the production of
39 scaffolds for medical use [12]. It is of interest, therefore, the understanding of how the initial material
40 properties and the process conditions interact with each other to determine the degree of
41 sintering/melting of the raw material powder, as well as the sintered material mechanical properties.
42 Most of the relevant investigations in the literature show that laser sintering is a very complicated
43 process [13], with several parameters influencing the densification mechanism and affecting the final
44 microstructural features.

45 Most of the current applications of SLS and SLM make use of metal [14,15] or polymeric powders
46 [16–19]. Most recent developments of these technologies have allowed to create functional objects
47 directly from metal powders without using either post-processing operations or binders (low melting
48 additives) [20]. A significant further spread of the application of these technologies can derive from
49 the use of ceramic materials [21]. The main issue with ceramic materials processing is the quality

50 improvement of the manufactured products [22]. In particular, size accuracy, strength and surface
51 finish of sintered parts are significantly affected by the type of material and by the process parameters
52 [23,24]. In fact, laser sintered artefacts are generally fragile and characterized by low density and
53 poor surface finish. Fragility of SLS/SLM artefacts is mainly due to high residual stresses in the
54 internal structure of the objects, resulting from the high temperature spatial non-homogeneity
55 determined by the process. In general, in all sintering processes the fragility of the artefact depends
56 on the strength of the interparticle solid necks resulting after solidification of molten contact zones
57 between particles [25]. This strength depends on the process parameters and mainly on the operating
58 temperature and pressure during the sintering process. For example, Luan et al. [26] found that the
59 neck tensile strength of sintered ash particles increases with the bonding time and pressure. Additional
60 influencing process parameters of SLS and SLM are given by the laser power and scan rate, the
61 scanning line spacing, the thickness of powder layers, the scanning geometry, the working
62 atmosphere, the powder bed temperature. Moreover, powder properties like particle material, particle
63 size distribution and particle shape play an important role on the final artifact characteristics as well
64 [27, 28]. The mechanical strength characterization of the artefacts includes the use of compression
65 [23] or bending [21] tests. In particular this latter can be applied to single layer objects, while the
66 former requires well defined 3D structures. In any case, destructive tests can be used to understand the
67 mechanical properties of the sintered artefacts [29]

68 Another important issue of the SLS process is the object shrinkage. This is influenced by material
69 and process parameters as well as by the geometry of the manufactured object. The object size and
70 shape inaccuracy due to shrinkage is generally overcome by calibrating the process and compensating
71 for it in the design of the initial powder bed specimen. Further improvements in the accuracy of the
72 manufactured object can be achieved by appropriate studies carried out to appreciate the sources of
73 the object size deviations, in order to adjust the calibration procedure to the process conditions [30].

74

75 The aim of the present work is to study the effects of some process parameters and powder properties
76 on the dimensions, the weight and the mechanical properties of the final sintered artefacts. In
77 particular, the study is focused on the effect of the laser power and the average particle size of ceramic
78 powders on the tensile strength the sintered materials. The application of a micromechanical model
79 allows to derive the interparticle neck size from the measured tensile strength.

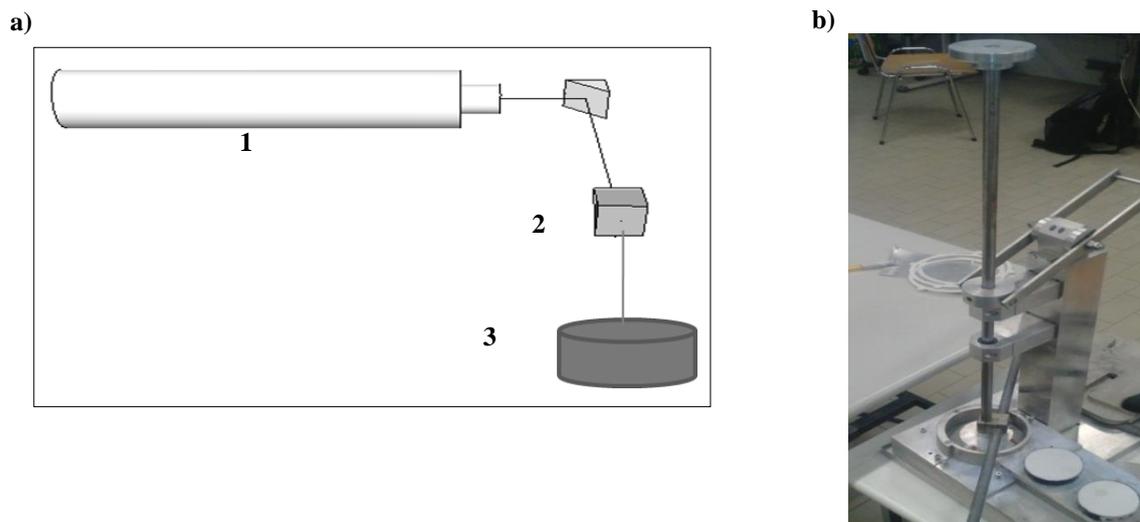
80 **2. Experimental methods**

81 *2.1 Apparatus*

82 The apparatus used in this work is reported in Figure 1a. It mainly consists of a laser tube and a system
83 to collimate the laser beam in order to hit in a precise position of the powder surface. In particular,
84 the CO₂ laser tube emits beams at a wavelength of 10.6 μm with a spot size of 100 micron diameter
85 and a maximum power of 40 W. The collimating device is made of a couple of inclined mirrors
86 moving on the x-y plane as shown in the figure. This system allows to change the laser scan velocity
87 in a range between 1 mm s⁻¹ and 100 mm s⁻¹. The nominal laser power of 40 W is sufficiently high to
88 allow the melting of ceramic materials in powder form. Adjusting the laser power and its speed it is
89 possible, therefore, to control the process in order to obtain melting or sintering of the powder
90 material. In principle, three dimensional objects might be obtained with the addition of new material
91 to be sintered in subsequent layers. The procedure generally used in commercial processes includes
92 the lowering of the powder sample, the spread and possibly the consolidation of the new material on
93 top of it, to re-establish the original powder height at the laser focal plane. The amount of lowering
94 defines the thickness of the fresh powder layer, that has to be thin enough to be effectively sintered
95 by the laser beam in its whole thickness. However, only the results obtained by working on a single
96 layer of sintered powder will be reported in this paper.

97 A special apparatus was used to consolidate and level beds of powders before sintering. The apparatus
98 is shown in Figure 1b and it consists of a flat lid connected by means of a cylindrical rod to a loading

99 plate. The vertical rod can freely move in linear bearings which allow the rod to move vertically and
100 to rotate around its axis. The rod has a loading plate on the upper end. This apparatus allows the
101 application of a pre-defined stress to the material below the lid, by applying the correct load on the
102 upper plate, and contemporarily, the use of a twisting action, obtained by rotating the rod with the
103 help of a lever.



104 Figure 1. Experimental set-up. a) Apparatus for laser sintering: 1. laser tube; 2. scanning mirrors moving in x
105 and y directions; 3. laser beam; b). Apparatus to consolidate and level beds of powders before the sintering
106 process.

107

108 2.2 Preparation of powder specimens

109 Glass beads, made of a soda-limestone glass, taken from three batches each characterized by a wide
110 particle size distribution (PSD) (<50 μm ; 10 to 100 μm ; 100 to 200 μm) were sieved in order to obtain
111 samples characterized by narrow PSD. For each of the sieved samples, Table 1 reports the nominal
112 ranges of the particle size of the original batches, the undersize and oversize sieve meshes used to
113 produce each sample, the mean Sauter diameter, d_s , of the sample measured by means of a laser
114 scattering measurement device (Malvern Mastersizer 2000), as well as the bulk density of the sample.
115 The particle density was measured with a helium pycnometer and it is equal to 2500 kg m^{-3} .

116

117 Table 1. Properties of sieved samples of glass beads.

Original cut size range, μm	Sieving range, μm	$d_s[3,2]$ μm	$d[4,3]$ μm	d_{10} μm	d_{50} μm	d_{90} μm	ρ_b kg m^{-3}
<50	< 25	16	15	10	15	20	1368
<50	25-32	27	28	17	25	34	1404
10-100	32-63	45	49	31	47	63	1504
100-200	63-125	86	87	56	85	114	1524
100-200	125-180	160	162	102	157	207	1536
100-200	160-250	184	182	119	180	241	1540

118

119 2.3 Sintering procedure

120 Single layer squared specimens (10x10 mm) of sintered materials were obtained in the laser apparatus
 121 described above. Prior to each sintering test, cylindrical beds of powder samples were levelled and
 122 consolidated at 2.5 kPa and two 90° twists operated by means of the powder consolidation apparatus
 123 described above. The sintering of squared specimens was obtained by moving the laser on fifty one
 124 10 mm long parallel lines at a distance of 0.2 mm from each other.

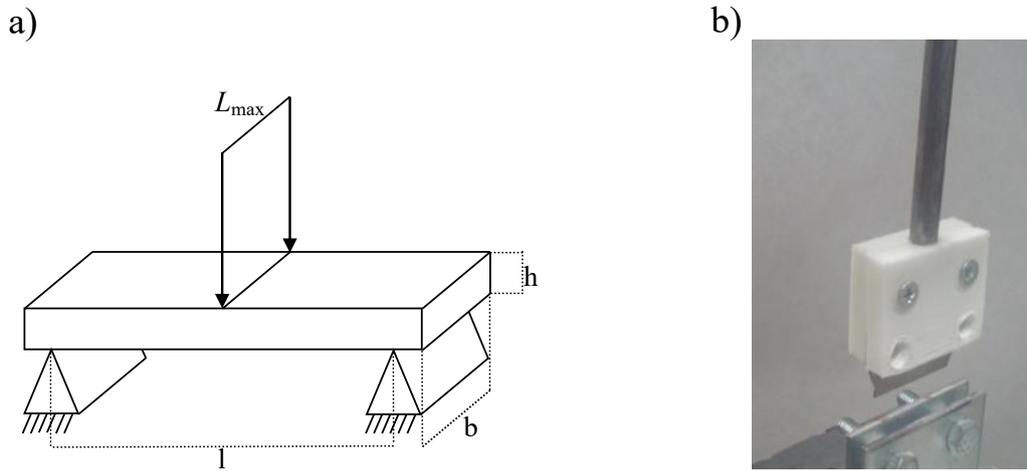
125 The laser scanning speed and the laser power were set according to two alternative combinations of
 126 values to provide the same scanning energy of 160 J m⁻¹. In particular, the first settings are given by
 127 a laser beam power of 8 W and a scan speed of 50 mm s⁻¹, while the second ones are given by a laser
 128 beam power of 16 W and a scan speed of 100 mm s⁻¹.

129 After the sintering process sintered specimens were gently extracted from the powder bed. Particles
 130 loosely adhering to the sintered specimen due to van der Waals forces were softly removed by means
 131 of a paint brush. The weight of sintered specimens was measured with an analytical balance (300
 132 Crystal).

133 2.4 Bending tests

134 Bending tests were carried out on the square shaped sintered specimens to relate the sintering process
 135 parameters to the tensile strength of the obtained artifacts. The tests consist in following the three
 136 point bending flexural method as sketched in Figure 2a. The instrument used to measure forces and

137 displacements is an Instron 5865 dynamometer equipped with a load cell of 2.5 N. Given the small
 138 size of the samples, a specifically designed tool, represented in Figure 2b, had to be used. The upper
 139 horizontal edges of two vertical plates kept at a known distance by two screws act as lower supports
 140 for the bending test. The blade hold at the tip of the load cell stem act as the upper bending tool.



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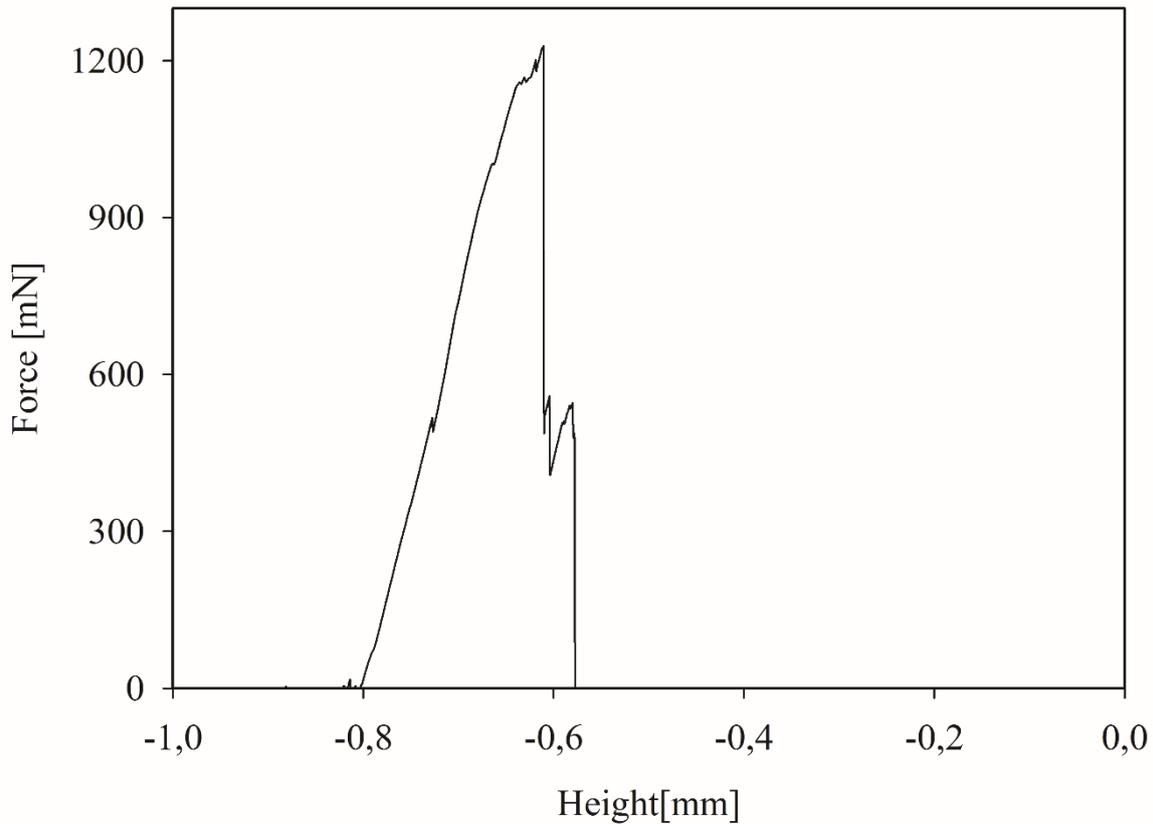
142 Figure 2. Three point bending test: a) Schematic; b) Bending tool.

143

144 A typical example of force displacement curve is reported in Figure 3. The x-axis reports the height
 145 of the edge of the upper bending tool above the holder, the y-axis reports the force measured by the
 146 upper bending tool. The layer thickness, t , can be derived from the difference between the height at
 147 which the bending tool starts to detect a finite force and the height of the supports. The peak force,
 148 L_{max} , identifies the sample breakage. Assuming a rectangular geometry of the sintered specimen the
 149 following equation can be used to relate the breakage force to the tensile strength, σ_b , of the sintered
 150 material:

$$151 \quad \sigma_b = \frac{M_b}{I} \frac{h}{2} = L_{max} \frac{3l}{2bt^2} \quad (1)$$

152 where M_b is the bending moment, I is the horizontal moment of inertia of the sample vertical cross
 153 section, h is the layer deflection, b is the specimen width, l is the distance between the supports.



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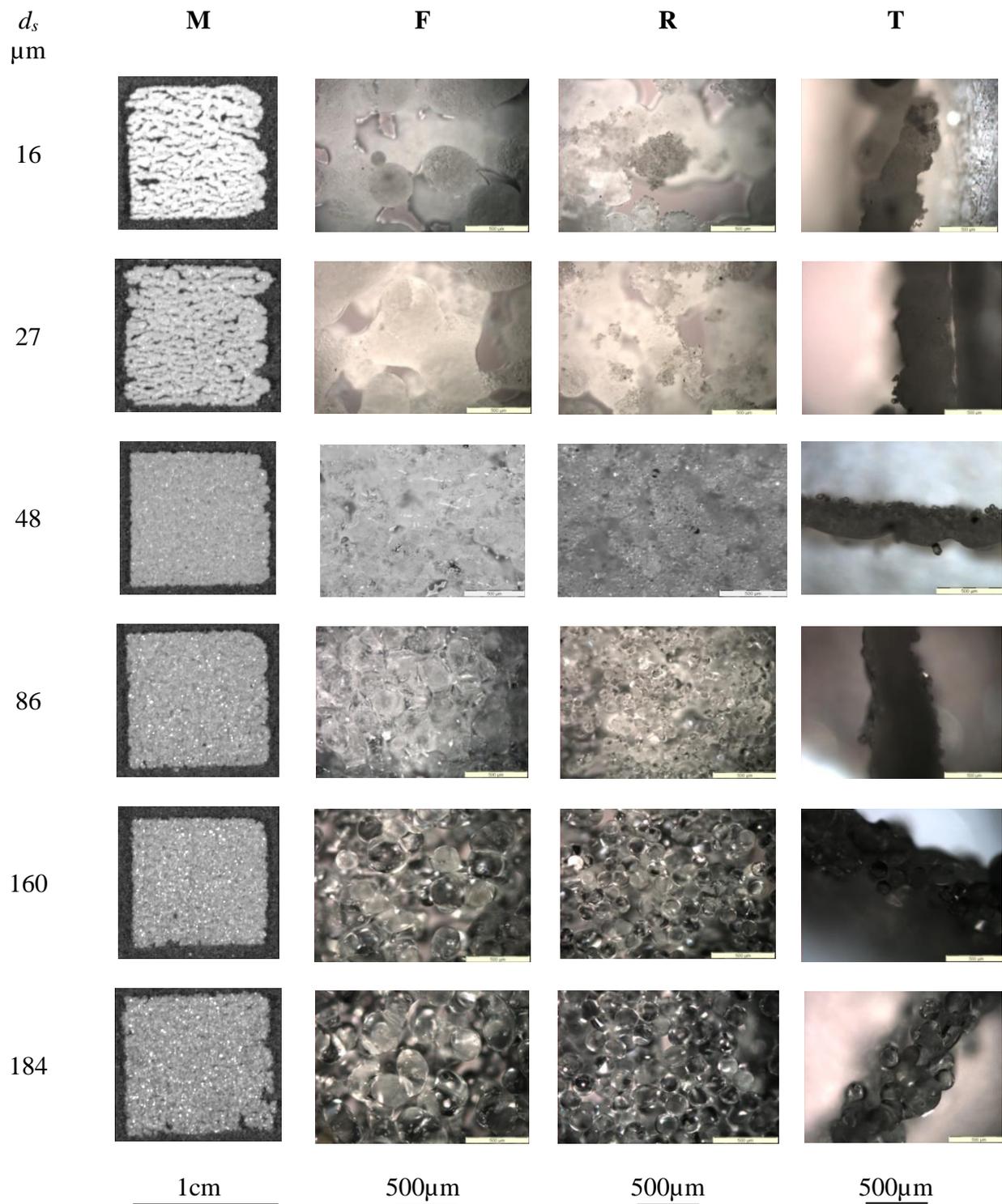
155 Figure 3. An example of a force vs displacement curve in a bending test.

156 3. Experimental results

157 3.1 Unimodal layers

158 Figure 4 reports macro and micro photographs of the sintered specimens using a laser beam of 8 W
 159 and a scan speed of 50 mm s^{-1} . In the Figure, organized in the form of a table, each row corresponds
 160 to a different mean diameter of the initial powder sample. The table starts with the finest sample
 161 ($d_s=16 \text{ }\mu\text{m}$) indicated with letter a) and ends with the coarsest sample ($d_s=184 \text{ }\mu\text{m}$) indicated with
 162 letter f)., Macro photographs of the whole $10 \times 10 \text{ mm}$ specimens are reported in column M of the
 163 Figure, while columns F, R and T report optical microscopy pictures (about 100x) of the sintered
 164 specimens, respectively the front (the laser sintering side) view, the back view and the lateral view.

165



166 Figure 4. Macro and micro photographs of 10x10 mm sintered specimens obtained from narrow PSD powder
 167 samples of different mean particle size d , by using a 8 W laser beam and a 50 mm s⁻¹ laser scanning speed. M
 168 macroscopic view; F, front view; R back view; T, lateral view.

169

170 Inspection of Figure 4 indicates that the particle size of the powder sample significantly affects the
 171 final morphology of the specimen by means of the extent of the molten zone with respect to the initial

172 particle dimensions. In particular, inspection of the front views (F) indicates that for finer powders
173 ($d_s=16\ \mu\text{m}$ and $d_s=27\ \mu\text{m}$) whole particles melted and, thus, the initial particle size and shape vanished
174 in the final solidified specimen made out of lumps. Instead, for coarser powders ($d_s=160\ \mu\text{m}$ and
175 $d_s=184\ \mu\text{m}$) melting was limited to interparticle contact zones and, thus, a final sintered specimen
176 with unchanged primary particle size and shape was obtained. Intermediate morphology was obtained
177 for powders with $d_s=48\ \mu\text{m}$ and $d_s=86\ \mu\text{m}$. Moreover, inspection of the microscopic back view (R)
178 of the specimens reveals the existence of regions of non-fused primary particles adhering to the fused
179 material. This evidence would suggest a non uniform laser energy transfer through the specimen
180 thickness. In particular, pictures of the specimens obtained from powder samples of $48\ \mu\text{m}$ and 86
181 μm seem to show melting on the front side and sintering on the back side. In the latter case a
182 continuous layer of sintered particles can be observed. Pictures of the specimens from powder
183 samples of $160\ \mu\text{m}$ and $184\ \mu\text{m}$ suggest that sintering is the predominating mechanism on the whole
184 layer thickness and that only some traces of melting are still visible on the front layer in form of
185 particles somewhat larger than the primary. This phenomenon might be reasonably attributed to the
186 coalescence of smaller fused particles during the process.

187 These results clearly indicate an effect of the particle size on the melting ability of the laser beam.
188 However, lateral views of the specimen (T) do not show significant differences in the thickness of
189 the specimen, suggesting that the penetration depth of the laser beam is comparable for all the tested
190 powders with different particle size. Further evidences on the effectiveness of laser penetration are
191 provided with the results of specimen weight, and of layer thickness, reported in Figure 5a and 5b
192 respectively. Inspection of Figure 5a reveals that the specimen weight tends to increase with
193 increasing particle size up to $d_s=160\ \mu\text{m}$, where melting is the prevailing phenomenon, and keeps
194 almost constant for larger particle sizes ($d_s=160\ \mu\text{m}$ and $d_s=184\ \mu\text{m}$), where sintering is occurring.
195 Layer thickness was measured in bending tests, according to the procedure explained in section 2.4.
196 Results reported in Figure 5b confirm that the specimen thickness does not change significantly with
197 the particle size in the range $27\text{-}160\ \mu\text{m}$. The smallest value of thickness was recorded for $d_s=16$

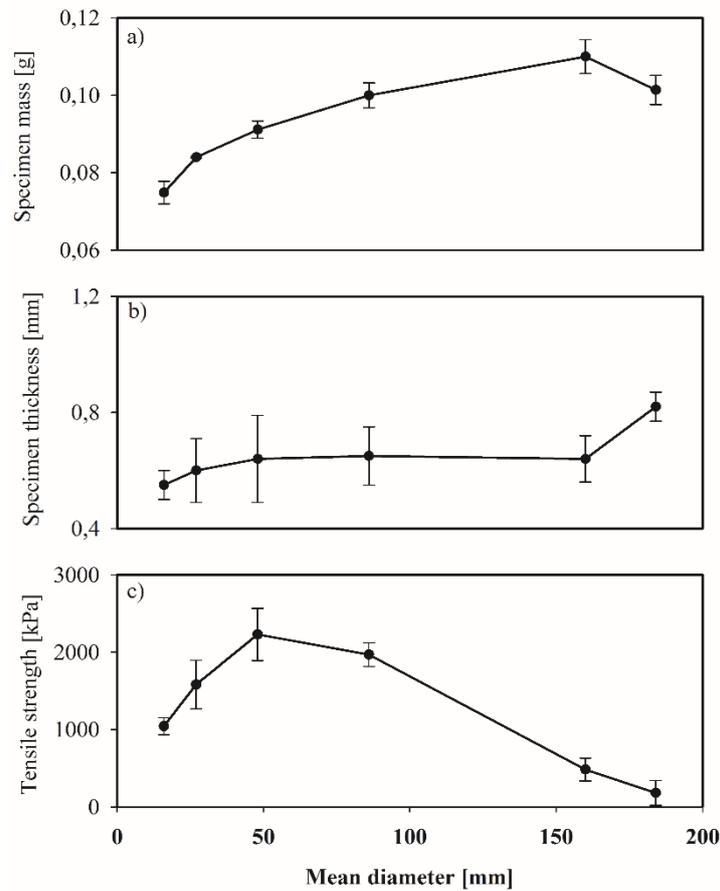
198 μm , while the largest value was measured for $d_s=184 \mu\text{m}$. Further inspection reveals that error bars
199 for specimens produced with finer powders ($d_s=16-48 \mu\text{m}$) are of the order of magnitude of the
200 particle size. This result might be explained with the adhesion of few particles on the specimen surface
201 due to van der Waals forces. In fact, van der Waal forces for rounded particles are comparable to the
202 particle weight up to diameters of the order of $100 \mu\text{m}$ (i.e page 124 [25] in Seville et al 1997).

203 As above mentioned, commenting results reported in Figure 4, the sintering process presents uneven
204 results along the thickness of the sintered specimen. In the bending test the specimen was positioned
205 on the supports keeping the orientation defined by the sintering process. This allowed measuring the
206 tensile strength of the material on its weaker side, that is the downward side. Tensile strength data
207 measured with this procedure are reported in Figure 5c. Inspection of the Figure reveals that the
208 tensile strength of the material presents a local maximum at $48 \mu\text{m}$. The tensile strength decreases
209 considerably for larger particle diameters with an almost tenfold reduction for the specimen sintered
210 from $184 \mu\text{m}$ particles with respect to specimen sintered from $27 \mu\text{m}$ particles. This reduction may
211 be attributed to two different causes: 1) the reduction of the number of sintered interparticle contact
212 points per unit volume with increasing particle size and 2) the reduced strength of single sintered
213 interparticle contacts due to different effects of the sintering process on particles of different sizes.

214 For simplicity, in the following the sintered contacts will be referred to as “necks”. Furthermore, as
215 discussed above, in the results, the obtained sintering of the finest samples, having mean particle
216 diameters between 16 and $27 \mu\text{m}$, probably correspond to process condition producing an almost
217 complete melting of the whole specimen thickness. Therefore, this can explain the change of trend
218 for the data relative to these points. In fact, the lower strength for finest particles might be attributed
219 to the lower bulk density of the initial powder, which brings to a looser final structure of the melted
220 material. Similar findings are reported by Tolochko [31] on experiments with SLM.

221 All the above considerations on the optical inspection of sintered specimens and on the measured
222 properties, seem to indicate that the sample characterized by an initial particle size of $48 \mu\text{m}$ marks a
223 transition between melting and sintering processes, given the laser power and the scanning speed used

224 in these experiments. All the results seem also to indicate that for specimen produced from powder
225 samples made of 86 μm particles and larger, particle sintering is the prevailing phenomenon.



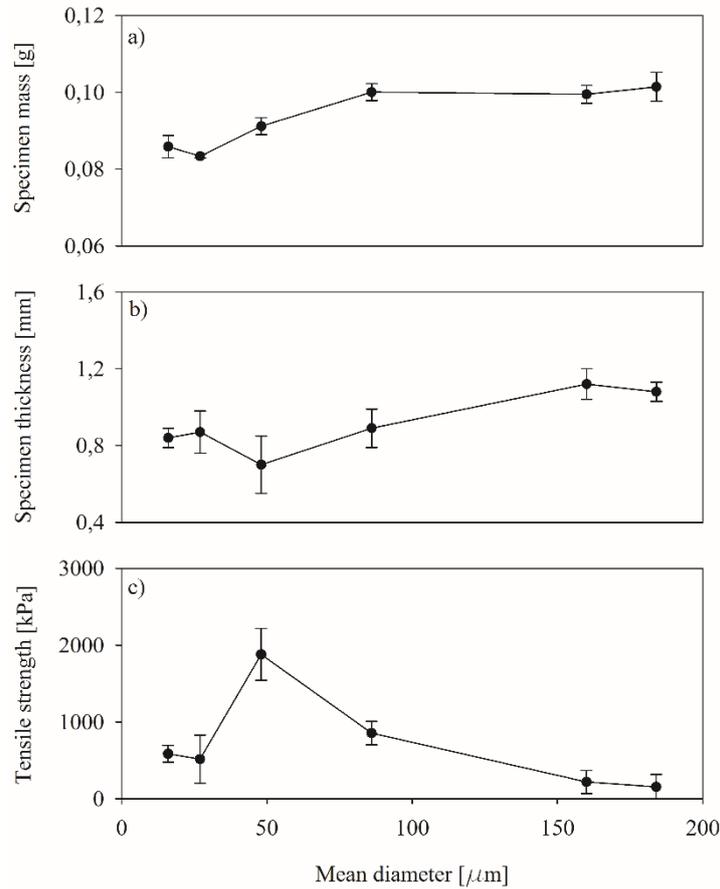
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227 Figure 5. Specimen properties as a function of the particle size of the initial unimodal powder sample:: a)
228 thickness; b) mass; c) tensile strength on the back side. 10x10 mm sintered specimens obtained by using a 8
229 W laser beam and a 50 mm s^{-1} laser scanning speed.

230

231 Specimen properties obtained starting from the same unimodal powder samples and using a double
232 laser scanning speed (laser power of 16W and scanning speed of 100 mm s^{-1}) are reported in Figure
233 6. Both the specimen mass and the specimen thickness for this case show similar values to those
234 obtained at low speed. Values in this case are just a bit higher for thickness. Furthermore, similarly
235 to the low scanning speed case, in general the specimen mass tends to increase with the particle size,
236 while the specimen thickness shows much more limited variations. Regarding the tensile strength, as
237 expected its order of magnitude in the high scanning speed case is the same to the one obtained in the
238 low speed case, however values in high scanning speed case are generally lower and, in particular for

239 the finest particles. These latter, differently from what observed at low scanning speed, show
240 specimen strength rather close to those of the coarsest particles.
241



242
243 Figure 6. Specimen properties as a function of the particle size of the initial unimodal powder sample: a)
244 thickness; b) mass; c) tensile strength on the back side. 10x10 mm sintered specimens obtained by using a 16
245 W laser beam and a 100 mm s⁻¹ laser scanning speed.
246

247 4. Discussion

248 4.2 Theoretical framework

249 The macroscopic properties of sintered artifacts will depend upon their microscopic structure, that is
250 the spatial density and the strength of the sintered contacts (necks). It might be argued that, in case of
251 particle sintering, that is in case of limited melting of the particles in the original powder system, the
252 spatial organization of particles can be assumed similar to the initial spatial organization of the

253 particles in the powders prepared for the laser sintering process. For powders a simple approach that
254 allows to relate mechanical properties of powders with the strength of the contact forces between
255 particles is the Rumpf [32] approach modified according to Molerus [33] suggestion, which allows
256 to estimate both the interparticle contact density and their role in the definition of powder strength.
257 The model was originally developed to relate van der Waals forces to powder strength in terms of
258 isostatic tensile stress and it was proven to apply satisfactorily also to uniaxial state of stresses. This
259 model however is sufficiently general to be used to correlate any kind of homogeneous distribution
260 of point contact force, compressive or tensile, on the particle surface to the powder stress,
261 compressive or tensile, locally acting on the powder. Therefore, applying this model at the failure
262 conditions in the following, this approach will be used to relate the tensile strength of the sintered
263 material with the tensile strength of single necks in it.

264 In the Rumpf and Molerus, model some simplifying assumptions are made, in particular it is assumed
265 that:

- 266 1. particles are organized in a randomly packed assembly and particles are spherical and
267 monodisperse;
- 268 2. the contact areas between particles are small enough in comparison with the particle surface and,
269 therefore, contact areas can be assumed as contact points;
- 270 3. the contact points are distributed over the particle spherical surface with equal probability;
- 271 4. the packing structure is isotropic;
- 272 5. the transmission of an isostatic state of stress, in which three equal principal stresses are assumed.

273 Starting from these hypotheses, Rumpf [32] and Molerus [33] derived the following equation relating
274 the isostatic stress, σ , with a mean isotropic contact force, F :

$$275 \quad \sigma = F \frac{k_0(1-\varepsilon)}{\pi d^2} \quad (2)$$

276 where d is the particle diameter, ε the particulate system void volume fraction (voidage) and k_0 the
277 particle coordination number. According to the correlation derived by Rumpf [32] from the

278 experimental data of Smith et al. [34] obtained for spherical particles and values of the particle
279 concentration between 0.55 and 0.65, the empirical relationship $k_0\varepsilon \approx \pi$ is assumed to estimate k :

$$280 \quad k_0 \approx \frac{\pi}{\varepsilon} \quad (3)$$

281 Considering the isostatic state of stress as the combination of three orthogonal uniaxial state of
282 stresses Molerus [33] extended the validity of equation (2) also to uniaxial state of stresses.

283 The simple model reported in equations (2) and (3) will be used to calculate the strength of the necks
284 in the sintered material of the specimen obtained from unimodal powder samples. The detail of this
285 procedure is given in the following.

286 Considering sintered specimen obtained by sintering materials with unimodal narrow PSD it is
287 possible to use equations (2) and (3) to estimate the strength of the sintered contact F_i for each average
288 particle diameter d_i starting from experimental values of the breakage modulus of the sintered material
289 σ_i

$$290 \quad F_i = \sigma_i \frac{d_i^2}{(1 - \varepsilon_i)} \varepsilon_i \quad (4)$$

291 the specific void volume fraction ε_i can be related to the experimental powder bulk density, ρ_{bi} , and
292 the particle density, ρ_p :

$$293 \quad \varepsilon_i = 1 - \frac{\rho_{bi}}{\rho_p} \quad (5)$$

294 The results of this calculation are provided in Table 2 for both the low and high scanning speed cases.
295 As it is possible to observe from this table, expectedly the particle coordination number k_{0i} is
296 increasing with the system bulk density, but in any case it is much smaller than the theoretical value
297 of 12, proper of the maximum packing density of equally sized spheres. The force of the sintered
298 contact F_i increases with the particle diameter, suggesting larger sintering necks for larger particles
299 and, rather strikingly, it is almost halved in the case of high scanning speed, in spite of the same

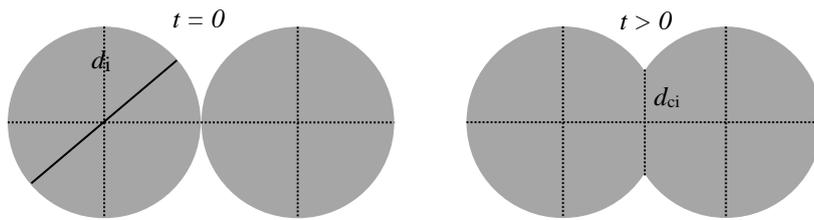
300 scanning energy used in the two tests. Assuming as a reference for the glass tensile strength σ_{tg} a
 301 value of 100 MPa it is possible to use the values of F_i to estimate the contact neck size, d_{ci} :

$$302 \quad d_{ci} = \sqrt{\frac{4F_i}{\sigma_{tg}\pi}} \quad (6)$$

303 As expected, the contact neck size increases with the particle diameter and, in any case, appears to be
 304 fully compatible with the particle diameters. Coherently with the smaller contact forces the neck
 305 diameters are smaller in the high scanning speed case.

306 In the selective laser sintering process the particle coalescence can be described by mechanisms
 307 similar to those in conventional sintering. A schematic of two particles in contact, undergoing
 308 sintering is shown in Fig.7. One of the earliest models to describe sintering in viscous regime was
 309 developed by Frenkel [35], who derived a rather simple expression to describe the time evolution of
 310 the neck radius, in which the neck diameter d_{ci} is normalized to the particle size d_i as follows:

$$311 \quad \left(\frac{d_{ci}}{d_i}\right)^2 = \frac{3}{\pi} \frac{\gamma(T_{gi})}{d_i \eta(T_{gi})} t \quad (7)$$



314
 315 Figure 7 Evolution in time of the dimensionless neck radius.
 316

317 In the Frenkel model, the sintering rate (or the neck size as a function of time) is dependent on γ , the
 318 surface tension, η , the viscosity of the melted material, and t , the sintering time. It states that sintering
 319 proceeds faster for materials with a higher surface tension, lower viscosity, and smaller particle
 320 radius. However, this model assumes that the particle radius remains nearly constant, and thus it can
 321 be valid only for the initial stage of the sintering, when the neck radius is much smaller than the
 322 particle radius, i.e. $d_{ci} \ll d_i$.

323 Equation (7) puts in evidence the fact that both the surface tension γ and the viscosity μ should be
 324 evaluated at the local temperature and, therefore, for its application it is necessary to have the right
 325 temperature dependence estimates [36]. Vice versa, equation (7) tells that if we know the neck size,
 326 the particle size, the sintering time and the temperature dependence of γ and μ it would be possible to
 327 use it to estimate an average temperature in the sintering process. Our glass beads are made of soda
 328 lime glass. For this type of materials, the temperature dependence equation is of the kind:

$$329 \quad \eta = A - \frac{B}{T_g - T_0} \quad (8)$$

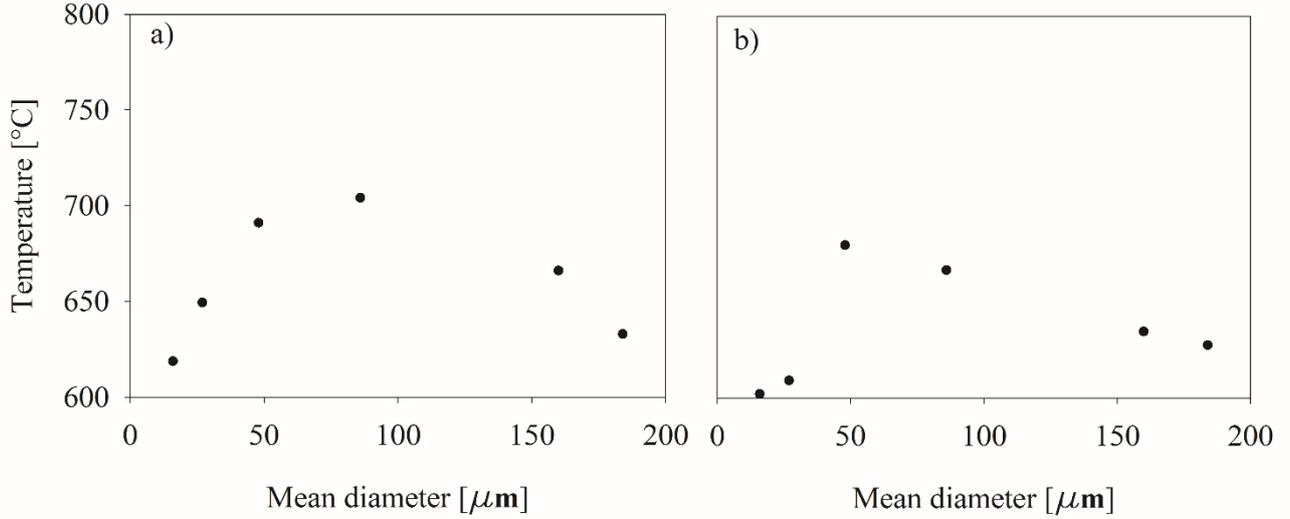
330 according to [37] and assuming a standard composition for the glass (in mass SiO₂ 71 % CaO 11%,
 331 Na₂O 18%) we adopt A=-2.54 Pa s, B=3810.7 Pa s °C⁻¹ and T₀=267.8 °C to calculate η in Pa·s and T
 332 in °C. Similarly, with reference to the surface tension, for this kind of glass [38] report values of s, r,
 333 t N·m⁻¹ at x, y and z °C, respectively. A second order polynomial was used to interpolate these values

$$334 \quad \gamma = aT_g^2 + bT_g + c \quad (9)$$

335 with $a= 3 \cdot 10^{-4}$ N m⁻¹ °C⁻², $b= -0.67$ N m⁻¹ °C⁻¹ and $c= 681.83$ N m⁻¹. As explained above, Equation
 336 (7) to (9) were solved with respect to the temperature given the particle size, the neck size and the
 337 sintering time t . In the lack of a better estimate this latter was set equal to the particle exposure time
 338 to laser t_{exp} that is equal to:

$$339 \quad t_{exp} = d_{sp} / V_{sc} \quad (10)$$

340 where d_{sp} is the laser spot size and V_{sc} the laser scanning speed. Figure 8 and Table 2 report the
 341 resulting averaged sintering temperatures T_g using the neck diameter reported in Table 2. It appears
 342 that the estimated sintering temperatures range for all systems between 600 and 700 °C. A maximum
 343 of temperature is found for intermediate particle size of the powder used, which is close to the values
 344 that provide the maximum strength of the sintered specimen.



345

346 Figure 8. Sintering temperatures calculated solving Eqs (7) to (9) using the experimental neck diameter
 347 reported in table 2, a) low speed, b) high speed.

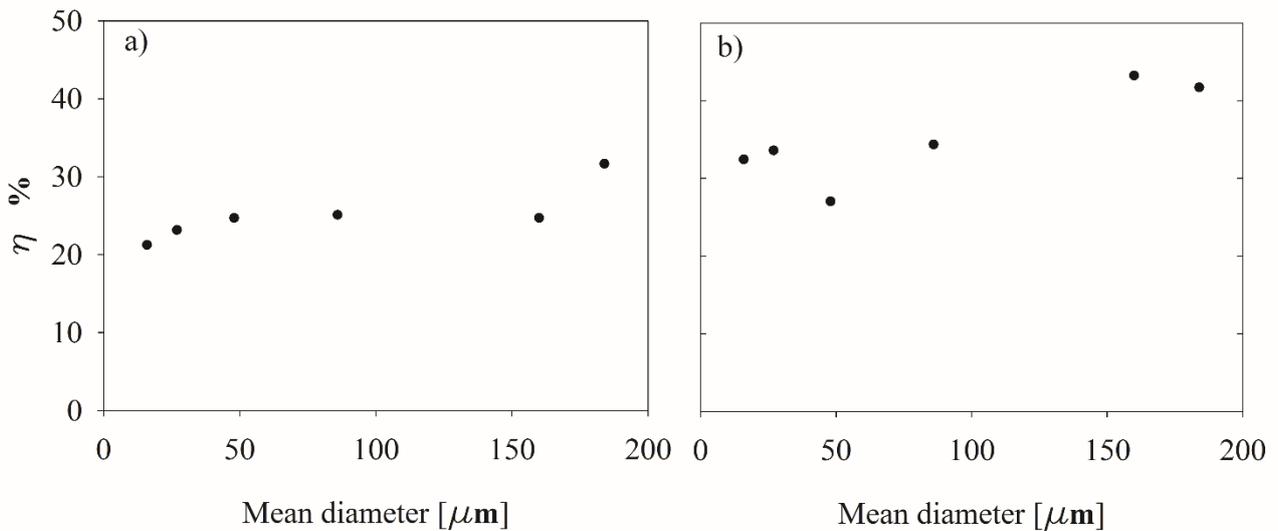
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349 With the underlying hypothesis that only the energy necessary for sintering would be desired, we can
 350 estimate the efficiency of the laser energy use. Given the particle exposure times, t_{exp} , given in Eq
 351 (10), it is possible to define the heat balance on the heated spot. It turns out that the efficiency of the
 352 laser energy use for the purpose of sintering, η_l , would be .

$$353 \quad \eta_l = \left[\frac{\pi d_{pi}^3}{6} \rho_p c_p (T_g - T_a) \left\{ \frac{\delta_l}{d_{pi}} \right\} \right] / \left(P_1 \frac{d_{pi}^2}{d_{sp}^2} t_{exp} \right) \quad (11)$$

354 In Eq (11), P_1 is the laser power, c_p is the specific heat of the particle material, T_g , is the sintering
 355 temperature, T_a the ambient temperature, δ_l the specimen thickness. In particular, the term in curly
 356 brackets in eq. (10) represents the average number of particles contained within the thickness of the
 357 sintered layer and accounts for the fact that the heating process does not involve a single particle layer
 358 only but the whole sintered layer thickness. The conditions tested are such that the product $P_1 \cdot V_{sc}$ is
 359 constant and that t_{exp} is equal to 2.0 ms for the low scanning speed case and it is 1.0 ms for the high
 360 scanning speed case. Values of η_l , reported in Table 2, indicate a general low value of the efficiency
 361 of the laser energy use. Higher values are found for large particles for which only sintering is observed
 362 and therefore energy is not used for particle the melting, produced by heating the particles above the

363 nominal temperature value of T_g , appearing in Eq. (11). Another source of low efficiency values of
 364 the laser energy use is likely to be the heat dispersion in the material around the laser spotlight. Higher
 365 efficiencies at higher scanning speeds seem to suggest more limited heat dispersion for shorter
 366 exposure times, which may also explain the different sintering activation of the contact points
 367 between large and small particles at the different scanning speeds.
 368



369
 370 Figure 9. Efficiency of the laser energy calculated solving Eqs (11), a) low speed, b) high speed.
 371

372 **5. Conclusions**

373 Microscopic observations on sintered materials and measures of mass, thickness and strength of
 374 single layer artifacts allow to characterize the effect of the main parameters in the laser sintering
 375 process.

376 The observation of the artifacts at the optical microscope allows to verify that the granular material
 377 directly subjected to irradiation of the laser has a high degree of fusion, while on the opposite side to
 378 that of the scan, the particles retain the original geometry. From experiments carried out on powder
 379 characterized by with narrow sized and unimodal particle size distribution it appears that an optimum
 380 in the sintered materials strength is obtained with particles sizes around 50 μm .

381 Table 2. Calculated sintered material properties according to eqs (4) to (6) at low scanning speed (laser power
 382 of 8W and scanning speed of 50 mm s⁻¹) and high scanning speed (laser power of 16W and scanning speed of
 383 100 mm s⁻¹).

Powder properties				Low scanning speed					High scanning speed				
d_{pi}	ρ_{bi}	ϕ_i	k	σ_i	F_{ci}	d_{ci}	η_l	T_g	σ_i	F_{ci}	d_{ci}	η_l	T_g
μm	kg m^{-3}	-	-	kPa	mN	μm	%	°C	kPa	mN	μm	%	°C
16	1370	0,547	6,85	1040	0,224	1,7	21	619	586	0,126	1,3	32	602
27	1410	0,562	7,07	1580	0,913	3,4	23	649	515	0,297	1,9	34	609
48	1480	0,58	7,39	2230	4,08	7,2	25	691	1880	3,17	6,4	27	680
86	1500	0,602	7,78	1970	9,77	11,2	25	704	855	4,24	7,4	34	666
160	1540	0,614	8,04	484	6,21	10	25	666	218	3,55	6,7	43	634
184	1540	0,616	8,07	182	3,89	7	32	633	155	3,32	6,5	42	628

384
 385 The model by Rumpf and Molerus, generally used to correlate the bulk properties of the powders to
 386 the interparticle forces, has been employed for the interpretation of the results in terms of material
 387 strength. In particular, we have estimated the change of the neck size of particles by changing the
 388 particle size and the laser scanning speed. This data coupled with the Frenkel model allowed an
 389 estimate of the sintering temperature for the different experimental conditions tested. The common
 390 outcome is that the optimum strength corresponds also to the largest neck size and the highest
 391 sintering temperature. Furthermore the critical analysis of results analysis, seems also to suggest that
 392 the main difference occurring in the sintering process between low and high scanning speeds at equal
 393 released energy in the sintering process is the higher effectiveness of laser beam at high speed, most
 394 likely for the more limited effects of energy losses from the laser sintering spot.

395 Nomenclature

- 396 c_p specific heat, W kg⁻¹ K⁻¹
 397 d_c , sintered contact diameter, m

398	d_p ,	particle diameter, m
399	d_s ,	mean sauter diameter, m
400	d_{sp} ,	size of the laser spot, m
401	f ,	finer fraction, -
402	F_c ,	interparticle force, N
403	k ,	particle coordination number, -
404	P_l	laser power, W
405	T_a	ambient temperature, K
406	t_{exp}	particle exposure time, s
407	T_g	glass transition temperature, K
408	t_{Hi}	particle heating time, s
409	V_{sc}	laser scanning speed, m s ⁻¹
410	W_{LF}	Effectiveness of contact between large and fine particles, -
411	<i>Greek</i>	
412	δ_l ,	specimen thickness, m
413	ρ_b ,	powder bulk density, kg m ⁻³
414	ρ_p ,	particle density, kg m ⁻³
415	σ ,	tensile strength, N m ⁻²
416	σ_{tg} ,	glass tensile strength, N m ⁻²
417	ϕ ,	particle volume fraction, -
418	<i>Subscripts</i>	
419	i,	referred to a single unimodal narrow PSD sample
420	b,	referred to a bimodal sample
421	F,	referred to the fine fraction

- 422 L, referred to the coarse fraction
423 LF, referred to the large particles in contact with fines

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